

# Statistical methods for variability studies

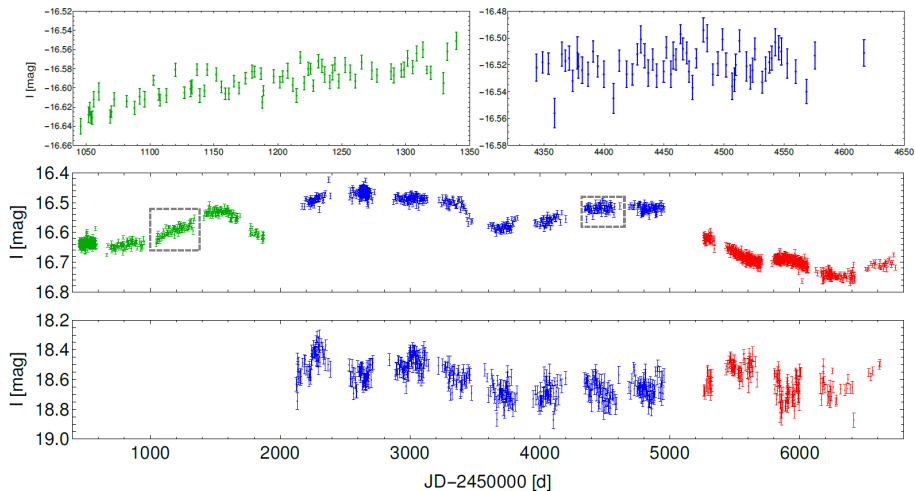
Mariusz Tarnopolski

Institute of Astronomy  
Nicolaus Copernicus University  
Toruń, Poland

WE-Heraeus Seminar  
Kraków, 7–10.11.2022

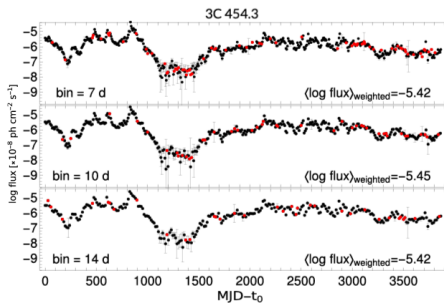
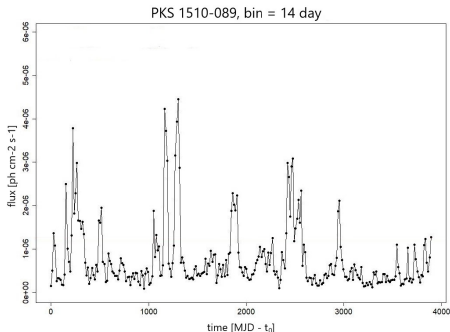


# Optical (OGLE) light curves of FSRQ and BL Lac candidates behind Magellanic Clouds:



Żywucka et al. (2018, 2020)

## Gamma-ray light curves of Fermi-LAT blazars:



# Power spectral density (PSD)

- 1 power law (PL):

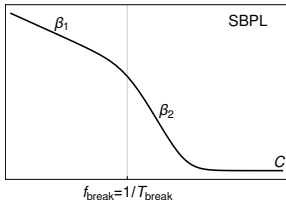
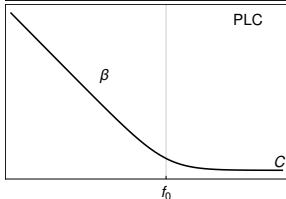
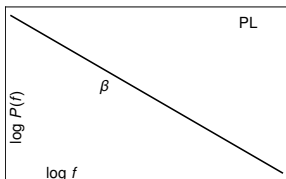
$$P(f) = \frac{P_{\text{norm}}}{f^\beta}$$

- 2 PL plus Poisson noise (PLC):

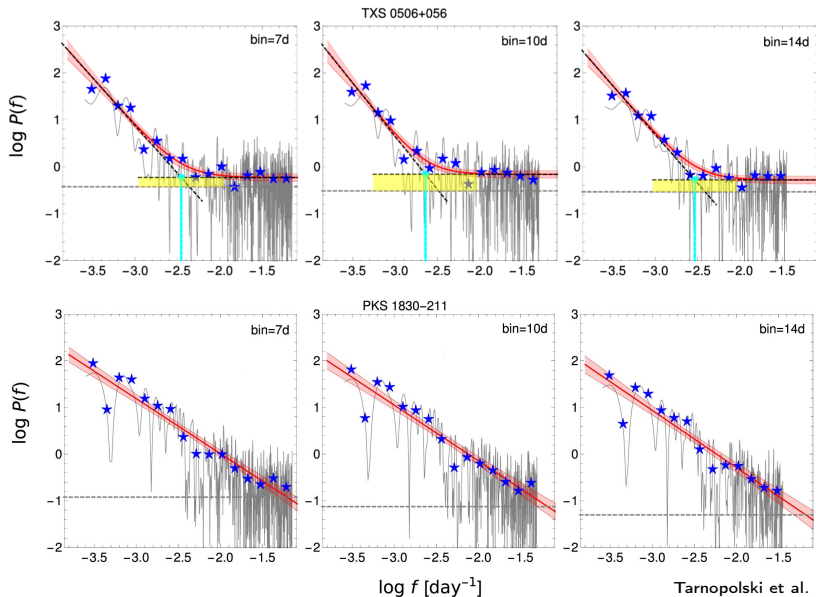
$$P(f) = \frac{P_{\text{norm}}}{f^\beta} + C$$

- 3 smoothly broken PL (SBPL) plus Poisson noise:

$$P(f) = \frac{P_{\text{norm}} f^{-\beta_1}}{1 + \left(\frac{f}{f_{\text{break}}}\right)^{\beta_2 - \beta_1}} + C$$

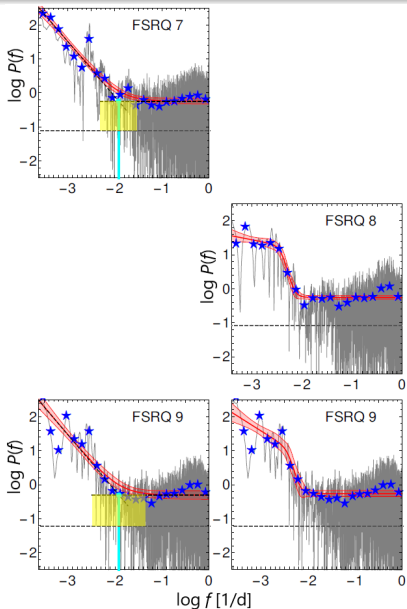


# Power spectral density (PSD) — Fermi-LAT



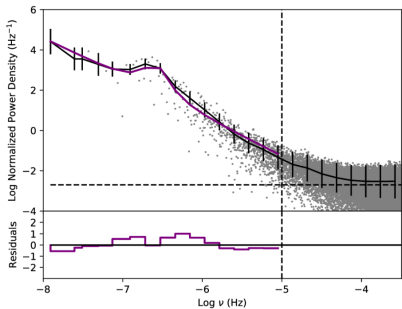
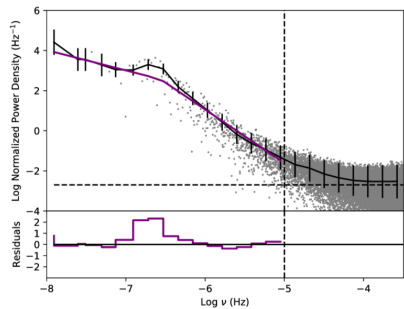
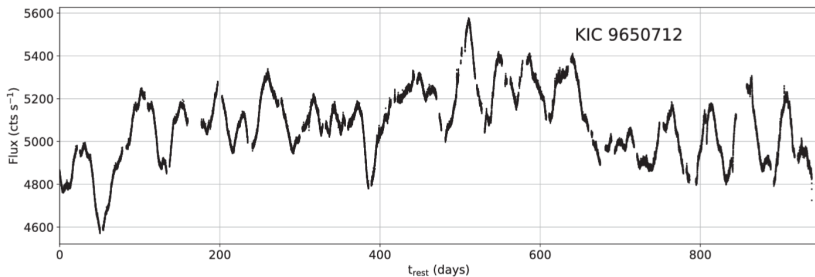
Tarnopolski et al. (2020)

# Power spectral density (PSD) — OGLE



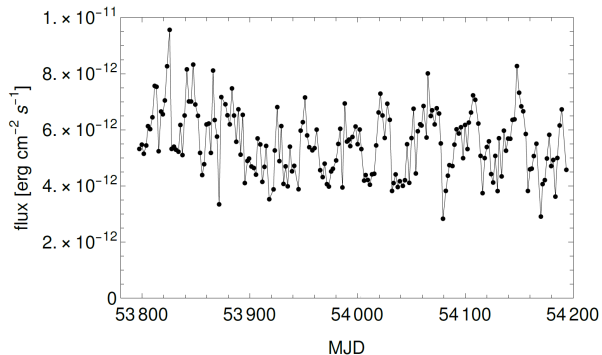
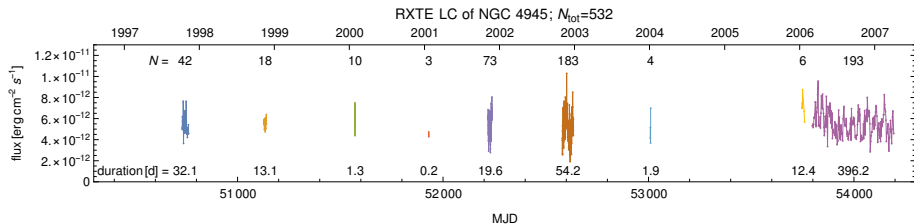
Żywucka et al. (2020)

# Quasiperiodic oscillations (QPOs) — Kepler



K.L. Smith et al. (2018)

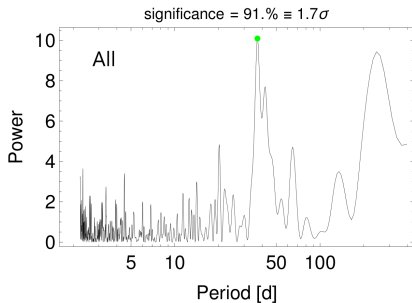
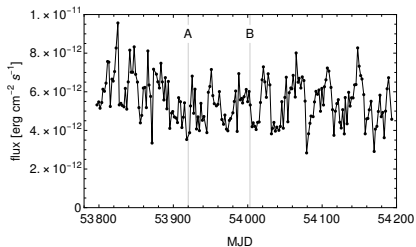
# Quasiperiodic oscillations (QPOs) — RXTE



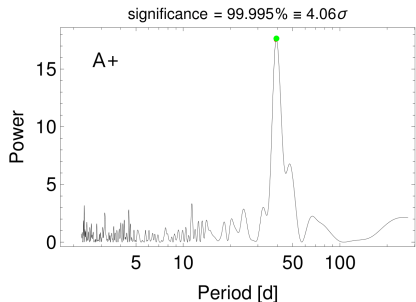
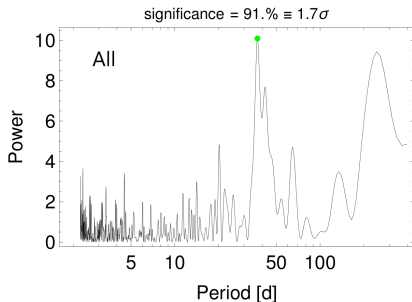
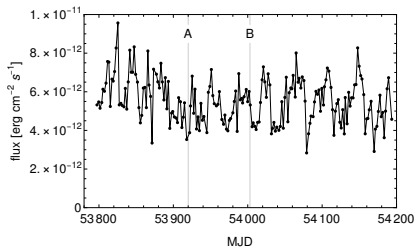
inspired by E. Smith et al. (2020)



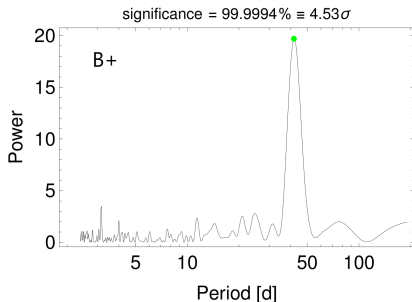
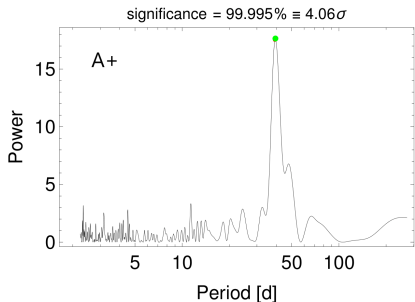
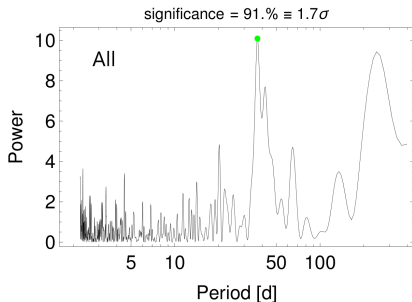
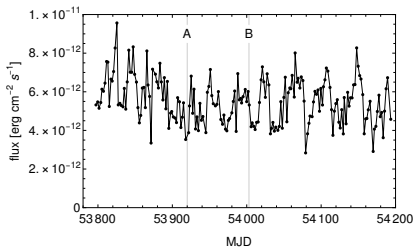
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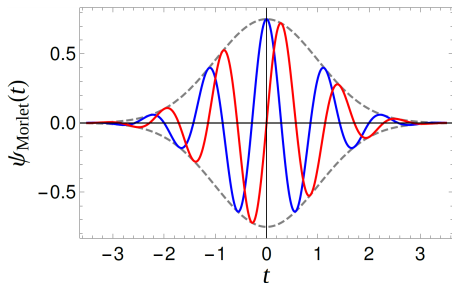
# Quasiperiodic oscillations (QPOs) — RXTE



MT & V. Marchenko (unpublished)

# Quasiperiodic oscillations (QPOs) via wavelets

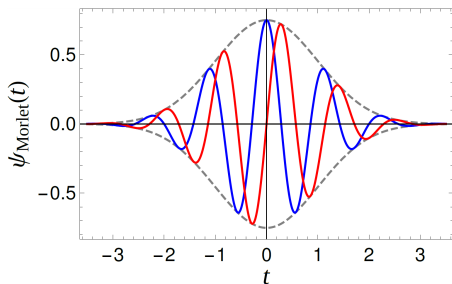
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A (mother) wavelet  $\psi(t)$  is a short, temporally and spectrally localized oscillation. Child wavelets form a basis ( $l$ —translation,  $s$ —scale):

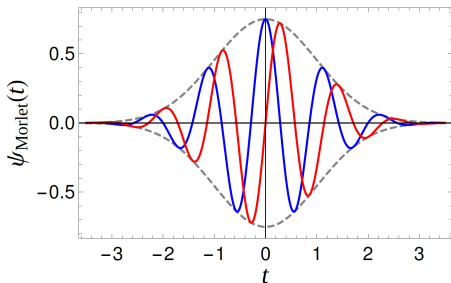
$$\psi_{s,l}(t) = \frac{1}{\sqrt{s}} \psi\left(\frac{t-l}{s}\right)$$



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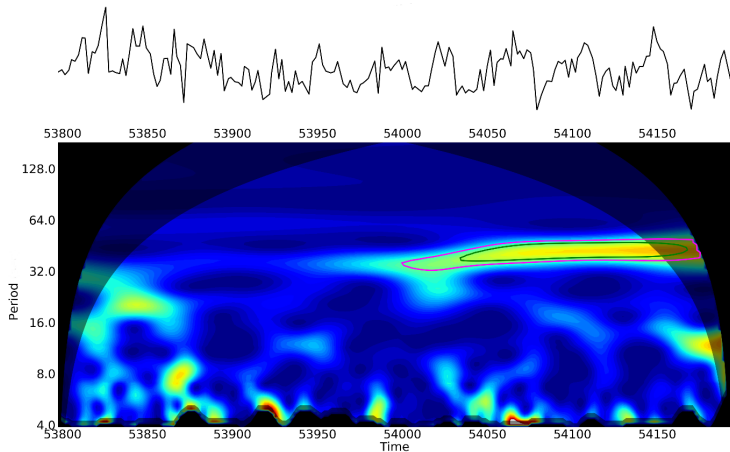


$$x(t) = \sum_{s,l} W(s,l) \psi_{s,l}(t)$$

$$W(s,l) = \int_t x(t) \psi_{s,l}^*(t) dt$$

$$P_{\text{wav}}(s,l) = |W(s,l)|^2$$

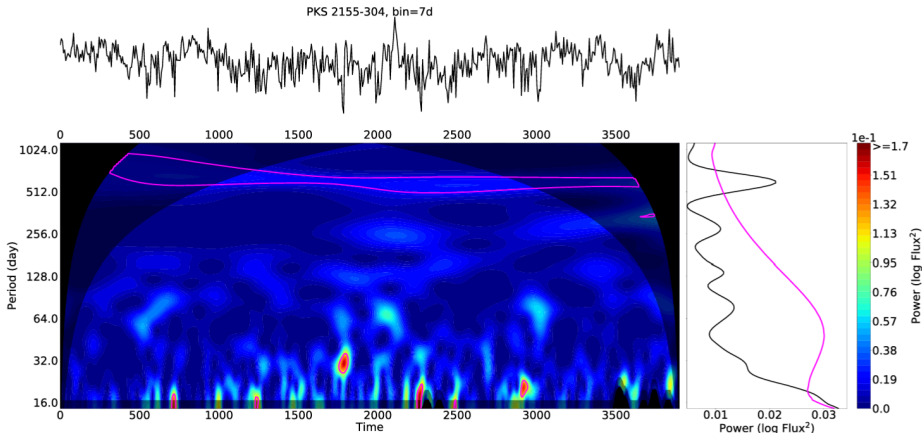
# Quasiperiodic oscillations (QPOs) — RXTE



Statistical significance via CARMA models

(cf. Sz. Kozłowski's talk & Kelly et al. 2014)

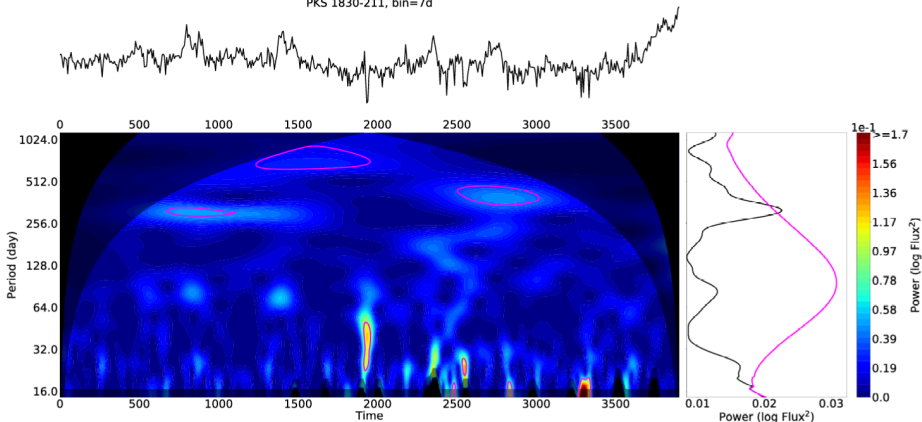
# Quasiperiodic oscillations (QPOs) — Fermi-LAT



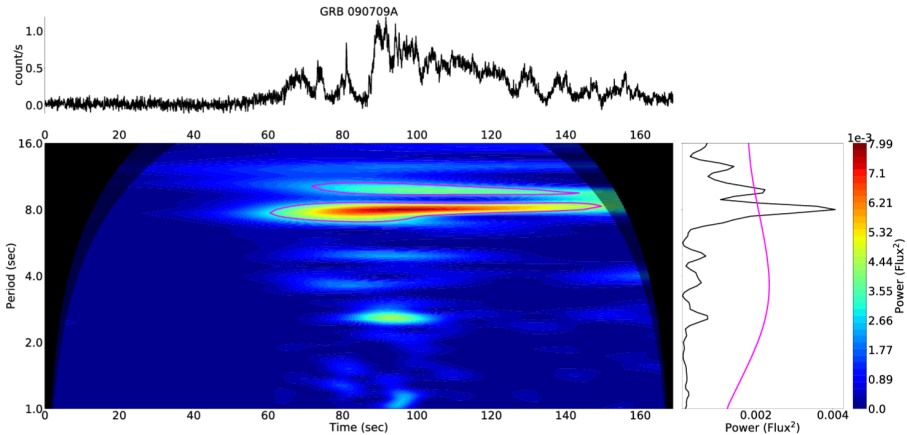


# Quasiperiodic oscillations (QPOs) — Fermi-LAT

PKS 1830-211, bin=7d



## Quasiperiodic oscillations (QPOs) — Swift-BAT GRBs

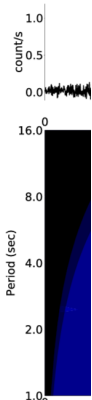


Tarnopolski & Marchenko (2021)

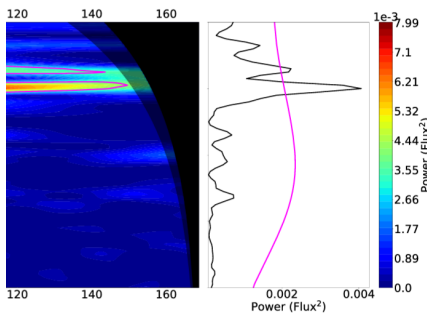
**Table 2**  
Identified QPOs

Number	GRB Name	Period (s)	Comment
6	GRB 200107B	$7.49 \pm 1.16; 11.40 \pm 1.57$	harmonics, 2 : 3
34	GRB 190821A	$8.20 \rightarrow 5.28$	up-chirp
75	GRB 190103B	$5.17 \pm 0.76$	constant
102	GRB 180823A	$19.12 \pm 3.19$	constant
122	GRB 180626A	$4.58 \pm 0.33; 5.70 \pm 0.54$	harmonics, 4 : 5
190	GRB 170823A	$2.96 \rightarrow 11.58$	down-chirp
212	GRB 170524B	$2.1 \rightarrow 2.8$	down-chirp
232	GRB 170205A	$6.86 \pm 0.80$	constant
250	GRB 161202A	$24.27 \rightarrow 16.25$	up-chirp
251	GRB 161129A	$3.83 \rightarrow 6.95$	up-chirp
252	GRB 161117B	$3.82 \pm 0.52$	constant
272	GRB 160824A	$3.05 \pm 0.56; 5.37 \pm 0.81;$ $9.43 \pm 1.51$	harmonics, 4 : 7 : 12 <sup>a</sup>
455	GRB 140730A	$12.32 \pm 1.95$	constant
462	GRB 140709B	$20.90 \pm 2.00;$ $41.54 \pm 4.30$	harmonics, 1 : 2
470	GRB 140619A	$8.87 \pm 0.99;$ $13.10 \pm 1.85;$ $32.34 \pm 3.86$	harmonics, 6 : 15 : 22 <sup>a</sup>
496	GRB 140323A	$5.49 \pm 0.98; 21.31 \pm 2.91$	harmonics, 1 : 4
551	GRB 130812A	$2.26 \pm 0.40$	constant
618	GRB 121209A	$9.89 \rightarrow 7.57$	up-chirp
622	GRB 121125A	$4.29 \pm 0.73; 8.48 \pm 1.00$	harmonics, 1 : 2
632	GRB 121014A	$16.70 \pm 1.87$	constant
701	GRB 120116A	$8.16 \pm 0.96$	constant
756	GRB 110422A	$5.46 \rightarrow 3.89$	up-chirp
777	GRB 110207A	$6.26 \pm 0.74$	constant
783	GRB 110107A	$5.48 \rightarrow 3.46$	up-chirp
805	GRB 100924A	$20.18 \rightarrow 5.14$	up-chirp
914	GRB 090709A	$8.02 \pm 0.67; 9.80 \pm 0.91$	harmonics, 4 : 5
945	GRB 090404	$10.94 \pm 0.86$	constant
963	GRB 090102	$7.64 \pm 1.07$	constant
1007	GRB 080810	$6.70 \pm 0.60; 9.15 \pm 0.85;$ $12.67 \pm 0.81$	harmonics, 2 : 3 : 4
1098	GRB 070911	$4.97 \pm 0.75; 16.50 \pm 2.08$	harmonics, 3 : 10
1127	GRB 070508	$2.14 \pm 0.26; 4.43 \pm 0.87$	harmonics, 1 : 2
1185	GRB 060906	$4.77 \pm 0.68$	constant
1324	GRB 050418	$14.70 \rightarrow 4.76$	up-chirp
1335	GRB 050306	$27.97 \pm 3.93$	constant

Quasiper



Swift-BAT GRBs



Tarnopolski & Marchenko (2021)

## Caveat emptor

Significance levels:

per cent	sigma
68%	$1\sigma$
90%	$1.64\sigma$
95%	$1.96\sigma$
99%	$2.58\sigma$
99.73%	$3\sigma$
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### Oops-Leon

- 6 GeV/ $c^2$  particle; significance 98% ( $\equiv 2.33\sigma$ ); denoted *Upsilon* ( $\Upsilon$ ).
- More data  $\Rightarrow$  spurious discovery.
- A 9.5 GeV/ $c^2$  particle discovered soon after at a  $5\sigma$  level—reused the name *Upsilon*.

Hom et al., PRL, 36, 1236 (1976); 39, 252 (1977)

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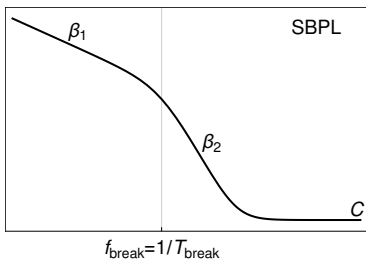
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# Black hole mass estimates



$$\begin{aligned} T_B &= 2\pi(r_{\text{ISCO}}^{3/2} + a_\star)(1+z) \frac{GM_{\text{BH}}}{c^3} \\ &= 0.359 m_9(1+z)f(a_\star) \text{ day} \end{aligned}$$

$$T_{\text{break}} \sim t_{\text{th}} \sim \alpha^{-1} t_K$$

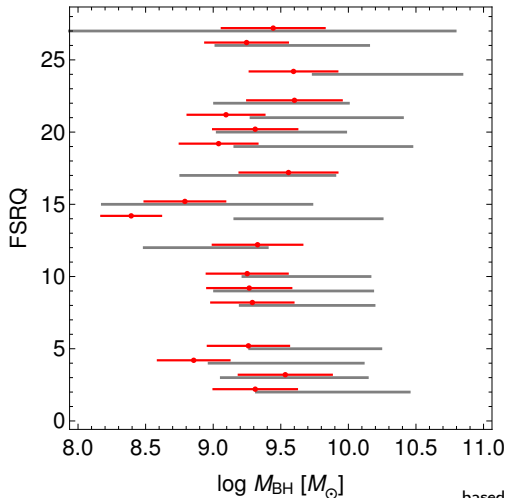
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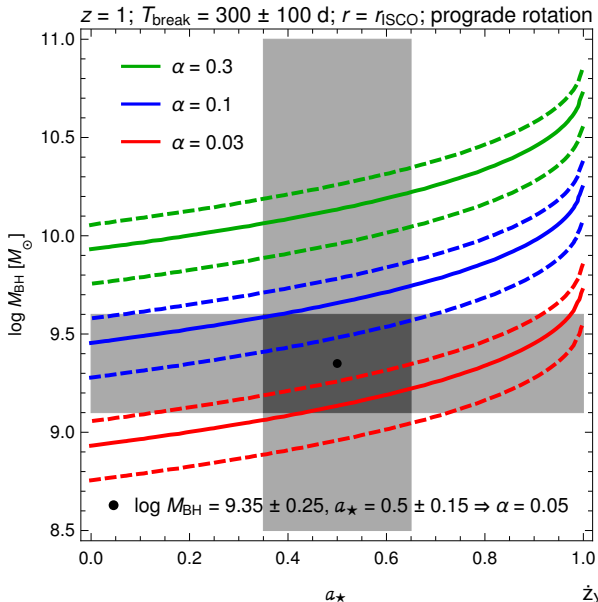
$$\log \frac{T_{\text{break}}}{1 \text{ day}} = A \log \frac{M_{\text{BH}}}{10^6 M_{\odot}} - B \log \frac{L_{\text{bol}}}{10^{44} \text{ erg s}^{-1}} + C$$

McHardy et al. (2006)

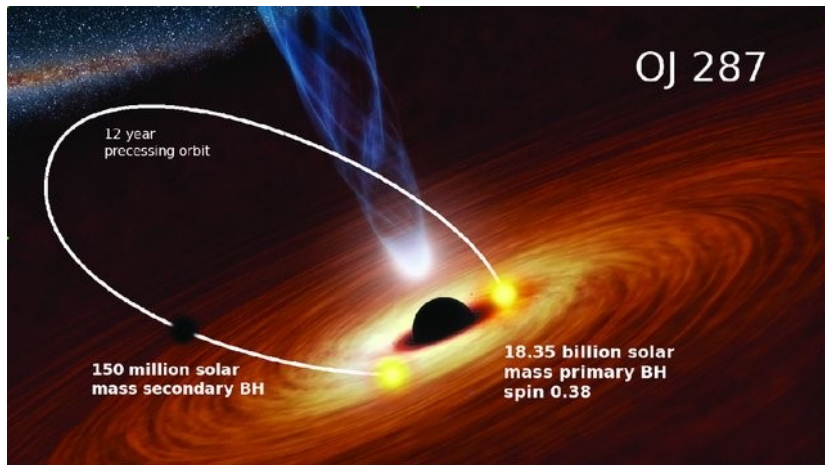


based on Żywucka et al. (2020)

# Viscosity estimates



# OJ 287 — binary SMBH



All that's periodic is not a binary black hole  
(Shakespeare, *travestied*, 1596)

- It has been speculated that even 10% of blazars can be a binary SMBH system
- They would contribute to the nHz GW background (cf. D. Champion's talk)
- Comparison with the pulsar timing array implies that binary SMBHs can constitute  $\lesssim 0.1\%$  blazars (Holgado et al. 2018)

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# QPOs — what else can cause them?

- Lense-Thirring eff. → disk precession —  $t_{LT} = \frac{8\pi GM}{c^3 a_*} \left(\frac{r}{R_S}\right)^3 \sim 0.1 \div 10 \text{ yrs}$
- Warped accretion disks
- Helical jet/magnetic fields; helical motion within the jet; twisting filaments in the jet etc.
- Perturbations in the accretion disk (matter or magnetic field densities, etc.) in the vicinity of the SMBH can propagate into the jet — however, e.g.:
  - variations in  $\dot{m}$  propagate through the disk with a time scale  $t_d \simeq \frac{1}{2\pi\alpha} \left(\frac{r}{H}\right)^2 t_K \sim 10^{2\div 3} \text{ yrs}$
  - but  $t_{th} \simeq \frac{1}{2\pi\alpha} t_K \sim 0.1 \div 10 \text{ yrs}$  (but  $a_* \neq 0$  in general)
- Jet's precession, even a slight one, changes the viewing angle — strong dependence of the Doppler factor  $\delta$  on the viewing angle → flux changes  $F_\nu = \delta^3 F'_\nu$  (also QPOs in polarization)
- Gravitational lensing (cf. D. Król's talk)
- Multitude of other scenarios — possibly a combination of several of them

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  - but  $t_{th} \simeq \frac{1}{2\pi\alpha} t_K \sim 0.1 \div 10 \text{ yrs}$  (but  $a_* \neq 0$  in general)
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- Gravitational lensing (cf. D. Król's talk)
- Multitude of other scenarios — possibly a combination of several of them

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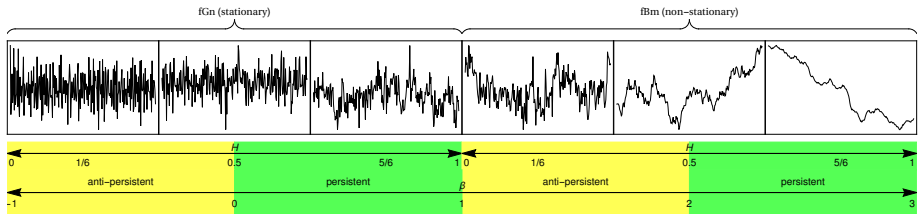
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# Hurst exponents

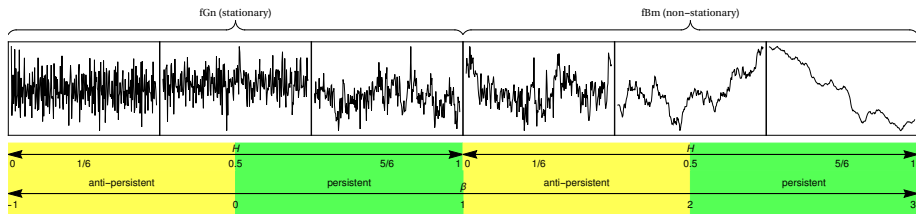
$$x(t) \doteq \lambda^{-H} x(\lambda t) \quad \rho_k \propto |k|^{-\delta} \equiv |k|^{-(2-2H)}$$





# Hurst exponents

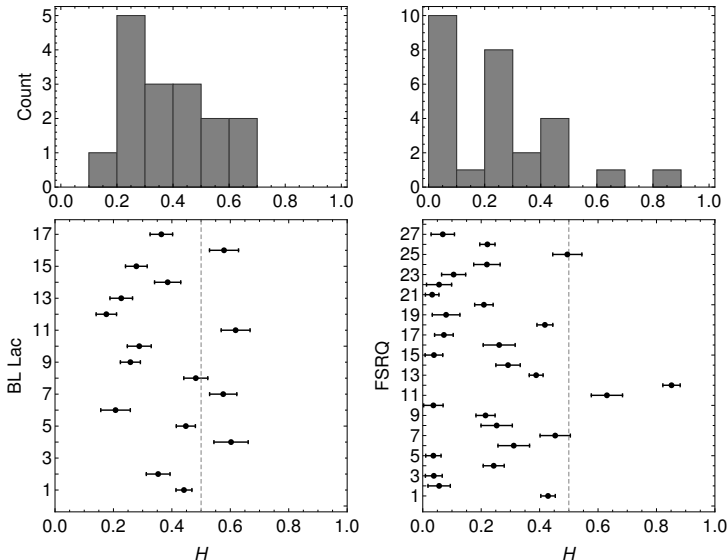
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The properties of  $H$ :

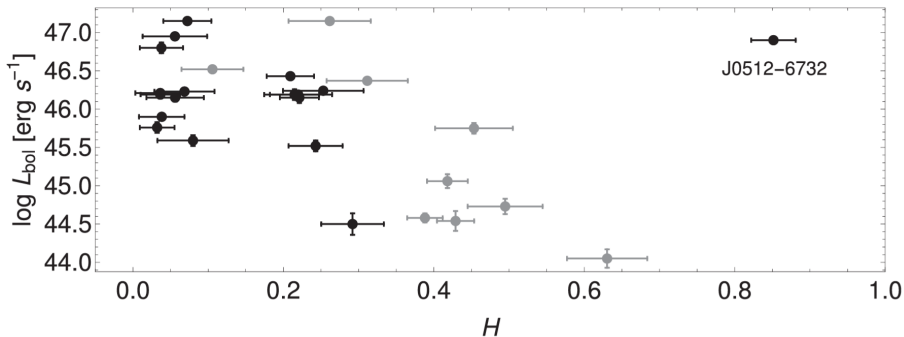
- ①  $0 < H < 1$ ,
- ②  $H = 1/2$  for an uncorrelated process,
- ③  $H > 1/2$  for a persistent (long-term memory, correlated) process,
- ④  $H < 1/2$  for an anti-persistent (short-term memory, anti-correlated) process.

# Hurst exponents — OGLE



Żywucka et al. (2020)

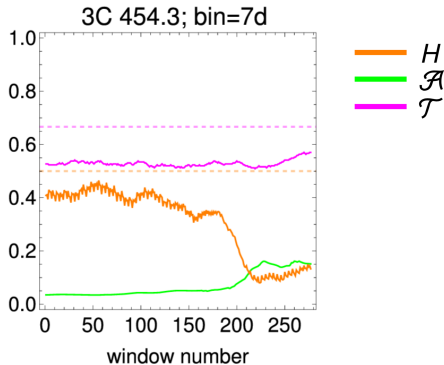
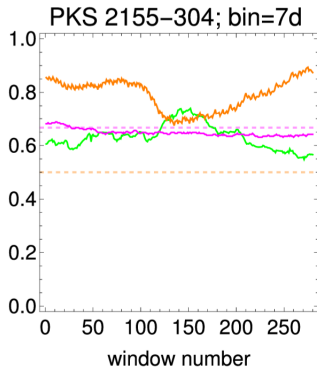
# Hurst exponents — OGLE



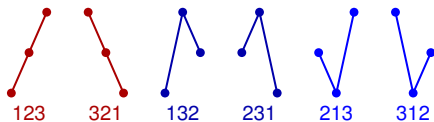
Black – SBPL; gray – PL.

$$r = -0.7$$

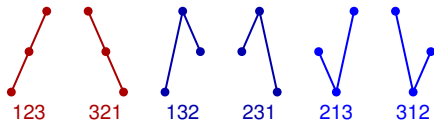
# Hurst exponents — Fermi-LAT



① Fraction of turning points,  $\mathcal{T}$ :



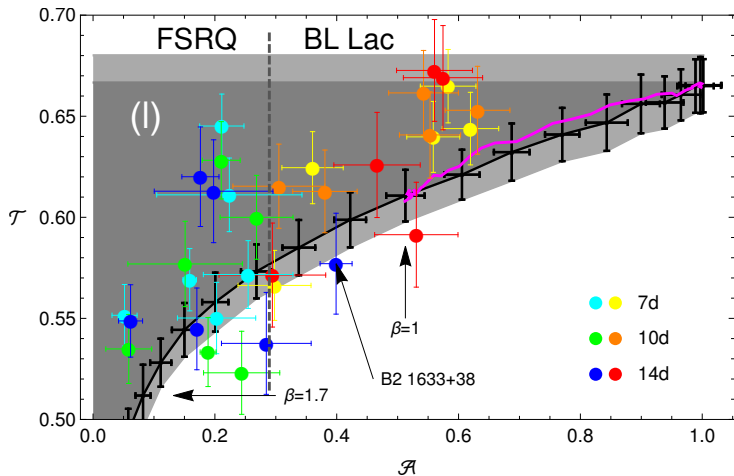
- ① Fraction of turning points,  $\mathcal{T}$ :



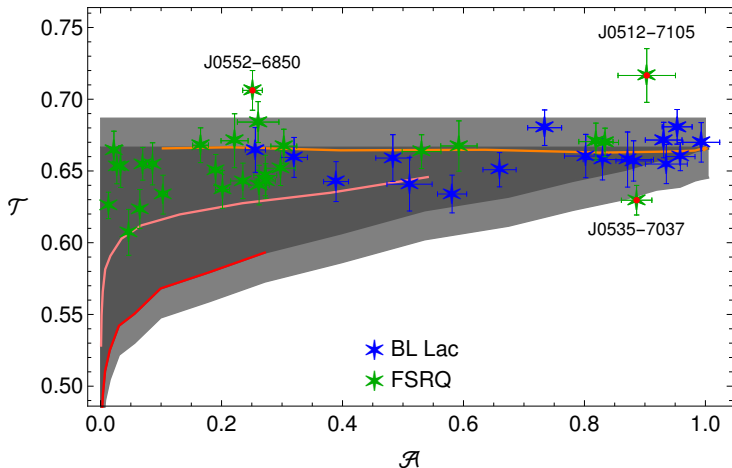
- ② Abbe value:

$$\mathcal{A} = \frac{\frac{1}{N-1} \sum_{k=1}^{N-1} (x_{k+1} - x_k)^2}{\frac{2}{N} \sum_{k=1}^N (x_k - \bar{x})^2} = \frac{1 \operatorname{var}(dX)}{2 \operatorname{var}(X)}$$

# $\mathcal{A} - \mathcal{T}$ plane — Fermi-LAT



# $\mathcal{A} - \mathcal{T}$ plane — OGLE



$$\langle \mathcal{A} \rangle_{\text{FSRQ}} = 0.29 \pm 0.05 \text{ and } \langle \mathcal{A} \rangle_{\text{BLLac}} = 0.71 \pm 0.06$$



- Wavelet scalograms, optimized for irregularly spaced data, and with **significance testing** — **crucial**
- Not every QPO indicates a binary SMBH
- Long, dense, high quality data → **more involved phenomenological stochastic models** — e.g., Hurst exponents capture some fine details about the autocorrelations
- $\mathcal{A} - \mathcal{T}$  plane — classification etc.
- **Connections with physical properties** — **important**

# Bibliography

- [1] Holgado, Sesana, Sandrinelli, et al., *MNRAS*, **481**, L74 (2018)
- [2] Kelly, Becker, Sobolewska, Siemiginowska, Uttley, *ApJ*, **788**, 33 (2014)
- [3] Lenoir & Crucifix, *NPGeo*, **25**, 175 (2018)
- [4] McHardy, Koerding, Knigge, Uttley, Fender, *Nat.*, **444**, 730 (2006)
- [5] Shakespeare, *The Merchant of Venice* (1596)
- [6] K.L. Smith, Mushotzky, Boyd, Wagoner, *ApJL*, **860**, L10 (2018)
- [7] E. Smith, Robles, Perlman, *ApJ*, **902**, 65 (2020)
- [8] Tarnopolski, *Physica A*, **461**, 662 (2016)
- [9] Tarnopolski, *PRE*, **100**, 062144 (2019)
- [10] Tarnopolski, Żywucka, Marchenko, Pascual-Granado, *ApJS*, **250**, 1 (2020)
- [11] Tarnopolski & Marchenko, *ApJ*, **911**, 20 (2021)
- [12] Żywucka, Goyal, Jamrozy, et al., *ApJ*, **867**, 131 (2018)
- [13] Żywucka, Tarnopolski, Böttcher, Stawarz, Marchenko, *ApJ*, **888**, 107 (2020)

☺ Thank you for your attention ☺